Symmetric Nuclear Matter (SNM)

Pure Neutron Matter (PNM)

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Conclusions

The properties of nuclear matter under the Bethe-Brueckner-Goldstone (BBG) Expansion

Verification of the convergence of the BBG expansion

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Overview				

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• The properties of nuclear matter

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- BHF Formalism
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3 Symmetric Nuclear Matter (SNM)

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Introduct	ion			

Properties of Nuclear Matter

- Saturation Density: $\rho = (0.17 \pm 0.01) \text{fm}^{-3}$
- Binding Energy: $\frac{E_0}{A} = (16 \pm 0.5) \text{MeV}$

The basic difficulty: traditional nucleon-nucleon forces feature a very strong repulsion at short distances.

We mainly employ **Bethe-Brueckner-Goldstone theory.**





¹ Z. H. Li et al. DOI: 10.1103/physrevc.74.047304

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Brueckner-Hartree-Fock Method

$$G(\rho,\beta;\omega) = v_{NN} + v_{NN} \sum_{k_1,k_2} \frac{|k_1k_2\rangle Q(1,2) \langle k_1k_2|}{\omega - \varepsilon(k_1) - \varepsilon(k_2)} G(\rho,\beta;\omega)$$
(1)

Nucleon-Nucleon interaction

$$V_{NN} = v_2 + V_3^{\text{eff}}$$

two-body interactions v_2 : AV18, CDBONN, N^3LO , N^4LO

• Single Particle Energy:

$$\varepsilon(k) = \frac{\hbar^2 k^2}{2m} + U(k)$$
 (2)

 \bullet "Auxiliary" potential, cont. choice: for all k

$$U(k) = \sum_{k' < k_F} \operatorname{Re} \left\langle kk' \big| G \big| kk' \right\rangle \tag{3}$$

• Gap (or standard) choice: $U(k > k_F) \equiv 0$



Fig 2: Different diagrams which contribute to the equation of state¹. $B/A = T + E_2 + E_3 + \dots$

(4)

¹ Song et al. DOI: 10.1103/physrevlett.81.1584

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Correlatior	1 Strength			

Definition

$$\kappa \equiv \sum_k n(k > k_F) / \sum_{k < k_F}$$

Coester et al. pointed out a <u>close connection</u> between the correlation parameter and the saturation properties of nuclear matter¹.

$$\kappa = \rho \int d^3 r \langle |\eta(r)|^2 \rangle_{S,T} = N \frac{V_{core}}{V} = \left(\frac{c}{d}\right)^3 \tag{5}$$

For a system with interaction range c and average particle distance d > c, the probability of finding a cluster n interacting particles is

$$\left(\frac{c}{d}\right)^{3n} = \kappa^n \ll 1 \tag{6}$$



Fig 3: Top Panel: Average depletion of the momentum distribution. Bottom Panel: Saturation curves of SNM in the BHF approximation².

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¹ Coester et al. DOI: 10.1103/PhysRevC.1.769

² Z.-H. Li and Schulze DOI: 10.1103/physrevc.94.024322

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Convergence of the hole-line expansion with modern nucleon-nucleon potentials

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We calculate the three-hole-line contributions to the binding energy of symmetric nuclear matter in the Brucckner-Bethe-Goldstone expansion using various modern nucleon-nucleon potentials of high precision. The relation with the correlation parameter $\kappa = \rho V_{core}$ is examined. In all cases the three-hole-line contributions turn out to be sufficiently small, but no satisfactory saturation is obtained. This means that three-nucleon forces are essential for all considered potentials.

DOI: 10.1103/PhysRevC.96.044309

Motivations

In the absence of a hard core, the very justification of the HLE becomes doubtful.

Explicitly compute the 3HL contributions.

Oteck the convergences of HLE with various modern NN potentials and different choices of aux. potentials.

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Binding I	Energy			



Fig 4: Saturation curves of symmetric nuclear matter for different NN potentials in the 2HL (left panel) and 2HL + 3HL (right panel) approximations with continuous (bold curves) or gap-choice (thin curves) s.p. potentials. The markers indicate the empirical saturation point.

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- Panel(a): From the density dependence of E₃ it becomes clear why no improvements: (Except N³LO500) all E₃ are decreasing function of density at ρ₀.
- 2 Panel(b) shows that for all potentials, the ratio $E_3/|E_2|$ remains well below 10% at any density, show again good convergence.
- Furthermore, no clear correspondence between the magnitudes of κ and E₃/|E₂| for in fact the total E₃ is the sum of the Bubble, Ring and Higher contributions.
- Panel(c): However, the contributions of E_{Bubble}/|E₂| are <u>completely consistent</u> with <u>κ</u> and orders according to the hard-core interaction strength.



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Defect Fu	nctions for			



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Fig 5: ${}^{1}S_{0}$ and ${}^{3}PF_{2}$ defect functions in PNM at different densities ($\rho_{0} = 0.17 \text{ fm}^{-3}$) for all potentials. The vertical dashed lines indicate the interparticle distance $d \equiv \rho^{-1/3}$.

r (fm)





Fig 6: r-space ${}^{1}S_{0}$ and ${}^{3}SD_{1}$ defect functions at different densities ($\rho_{0} = 0.17 \text{ fm}^{-3}$) for different potentials.

¹ Z.-H. Li and Schulze DOI: 10.1103/physrevc.94.024322





Fig 7: Energy per particle of PNM for different NN potentials in 2HL (left panel) and 2HL+3HL (right panel) approximation with continuous (thick curves) or gap-choice (thin curves) s.p. potentials. The N⁴LO550 calculations do not converge well for $\rho \gtrsim 0.5 \ {\rm fm}^{-3}$.

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Conclusio	ns			

HLE are well converged

- With all frequently used potentials, the 3HL contributions are sufficiently small.
- Results with gap- and continuous- choice show better agreement after including 3HL results.

For SNM, the empirical saturation properties of nuclear matter are not reproduced for any potentials which means that very strong nuclear three-body forces are required in order to achieve satisfactory saturation properties of nuclear matter.

The relation with the correlation parameter $\kappa = \rho V_{core}$ is examined.

We analyze the potentials in terms of the strength of their hard core.

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